

Ensemble Kalman filtering in the near-Gaussian setting

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The discrete-time filtering problem

Analysis of a simple Gaussian projection filter

Analysis of the ensemble Kalman filter in the near-Gaussian setting

Conclusions and perspectives

The discrete-time filtering problem

We consider the following stochastic dynamics and data model:

Dynamics and observations

Stochastic dynamics:	$u_{n+1} = \Psi(u_n) + \xi_n,$	$\xi_n \sim N(0, \Sigma),$
Data model:	$y_{n+1} = h(u_{n+1}) + \eta_{n+1},$	$\eta_{n+1} \sim N(0, \Gamma).$

Independence assumption:

$$u_0 \perp \xi_n \perp \eta_n$$

Initial state: $u_0 \sim N(m_0, C_0)$.

Notations:

- $\{u_n\}_{n \in [0,N]}$ is the unknown state in \mathbb{R}^d .
- $\{y_n\}_{n \in [\![1,N]\!]}$ are the observations in \mathbb{R}^K .
- $\Psi \colon \mathbf{R}^d \to \mathbf{R}^d$ and $h \colon \mathbf{R}^d \to \mathbf{R}^K$ are nonlinear operators.
- $Y_n = \{y_1^{\dagger}, \dots, y_n^{\dagger}\}$ is a given realization of the data up to time n.

Goal: Approximate sequentially the probability measure $\mathbf{P}(u_n|Y_n)$ for $n \in [\![1,N]\!]$.

Evolution of the filtering distribution

Update formula for the true filtering distribution

Let $\rho_n = \mathbf{P}(u_{n+1}|Y_{n+1})$. Then

 $\rho_{n+1} = L_n P \rho_n$

For simplicity, we assume that all the measures have densities.

• The map $P \colon \mathcal{P}(\mathbf{R}^d) \to \mathcal{P}(\mathbf{R}^d)$ is the prediction:

$$\frac{P}{\rho(u)} = \frac{1}{\sqrt{(2\pi)^d \det \Sigma}} \int \exp\left(-\frac{1}{2} |u - \Psi(v)|_{\Sigma}^2\right) \rho(v) \,\mathrm{d}v.$$

• The map $L_n: \mathcal{P}(\mathbf{R}^d) \to \mathcal{P}(\mathbf{R}^d)$ is the analysis (Bayes' theorem):

$$L_n \rho(u) = \frac{\exp\left(-\frac{1}{2} \left|y_{n+1}^{\dagger} - h(u)\right|_{\Gamma}^2\right) \rho(u)}{\int_{\mathbf{R}^d} \exp\left(-\frac{1}{2} \left|y_{n+1}^{\dagger} - h(U)\right|_{\Gamma}^2\right) \rho(U) \,\mathrm{d}U}$$

Schematically,

$$\mathbf{P}(u_n|Y_n) \xrightarrow{P} \mathbf{P}(u_{n+1}|Y_n) \xrightarrow{L_n} \mathbf{P}(u_{n+1}|Y_{n+1}).$$

The discrete-time filtering problem

Bootstrap particle filter

One of the simplest filtering methods is based on the iteration

$$\rho_{n+1}^B = L_n S^J P \rho_n^B.$$

where S^{J} is the sampling operator

$$S^J
ho := rac{1}{J} \sum_{j=1}^J \delta_{u^{(j)}}, \qquad u^{(j)} \sim
ho ext{ i.i.d.}$$

- The filtering distribution is approximated by a sum of Dirac masses;
- The convergence of this approach can be proved^{[1],[2]} in the metric

$$d_1(\mu, \nu) := \sup_{\|f\|_{\infty} \le 1} \mathbf{E} \sqrt{|\mu[f] - \nu[f]|^2}.$$

P. REBESCHINI and R. van HANDEL. Can local particle filters beat the curse of dimensionality? Ann. Appl. Probab., 2015.

^[2] D. SANZ-ALONSO, A. M. STUART, and A. TAEB. Inverse Problems and Data Assimilation with Connections to Machine Learning. arXiv preprint, 2018.

The bootstrap particle filter

- \blacksquare converges to the true filtering distribution in the limit $J \to \infty,$
- but tends to perform poorly for high-dimensional problems^[3].

In this talk, we study the ensemble Kalman filter,

- more robust in high-dimensional setting,
- but converges to the right limit as $J \to \infty$ only in the Gaussian setting.

Objective: obtain an error estimate away from the Gaussian setting.

^[3] P. BICKEL, B. LI, and T. BENGTSSON. Sharp failure rates for the bootstrap particle filter in high dimensions. In Pushing the limits of contemporary statistics: contributions in honor of Jayanta K. Ghosh, Inst. Math. Stat. (IMS) Collect. 2008.

Using stochastic filtering to solve inverse problems (1/2)

Inverse problem considered

Find an unknown parameter $oldsymbol{u} \in \mathbf{R}^d$ from data $y \in \mathbf{R}^K$ where

 $y = h(\mathbf{u}) + \eta,$

- h is the forward operator;
- η is observational noise.

Bayesian approach:

- Probability distribution on parameter: $u \sim \rho_0$, encoding **prior knowledge**;
- Probability distribution for the noise: $\eta \sim \nu$.

Bayesian **posterior distribution** given observation y^{\dagger} :

$$ho^y({m u}) \propto
ho_0({m u}) \,
uig(y^\dagger - h({m u})ig) = {\sf prior} imes {\sf likelihood}.$$

In the Gaussian case where $\nu = N(0, \Gamma)$,

$$\rho^{y}(\boldsymbol{u}) = \exp\left(-\left(\frac{1}{2}\left|y^{\dagger}-h(\boldsymbol{u})\right|_{\Gamma}^{2}\right)\right)\rho_{0}(\boldsymbol{u}).$$

Connection with stochastic filtering

The posterior distribution coincides with the filtering distribution at n = N for

$$u_{n+1} = u_n,$$

 $y_{n+1} = h(u_{n+1}) + \eta_{n+1}, \qquad \eta_{n+1} \sim \mathsf{N}(0, N\Gamma).$

with observations $y_n^{\dagger} = y^{\dagger}$ and initial state $u_0 \sim \rho_0$.

Indeed, in this case P is the identity map and so

$$\rho_{n+1} = L\rho_n = \frac{\exp\left(-\frac{1}{2}|y^{\dagger} - h(u)|_{N\Gamma}^2\right)\rho_n(u)}{\int_{\mathbf{R}^d} \exp\left(-\frac{1}{2}|y^{\dagger} - h(U)|_{N\Gamma}^2\right)\rho_n(U)\,\mathrm{d}U}$$
$$= \dots = \frac{\exp\left(-\frac{1}{2}|y^{\dagger} - h(u)|_{N\Gamma/n}^2\right)\rho_0(u)}{\int_{\mathbf{R}^d} \exp\left(-\frac{1}{2}|y^{\dagger} - h(U)|_{N\Gamma/n}^2\right)\rho_0(U)\,\mathrm{d}U}$$

 \rightsquigarrow Numerical methods for filtering are useful in the context of inverse problems.

Example: inference of the thermal conductivity in a plate



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Analysis of a simple Gaussian projection scheme

A simple (mean field) filter based on Gaussian projections

We consider the dynamics

$$\rho_{n+1}^G = GL_n P \rho_n^G, \qquad G\rho := \mathsf{N}\big(\mathcal{M}(\rho), \mathcal{C}(\rho)\big).$$

where

$$\mathcal{M}(\rho) = \int \theta \rho(\mathrm{d}\theta) \,, \qquad \mathcal{C}(\rho) = \int \big(\theta - \mathcal{M}(\rho)\big) \otimes \big(\theta - \mathcal{M}(\rho)\big) \rho(\mathrm{d}\theta) \,.$$

- If Ψ and h are linear, then $\rho_n^G = \rho_n$.
- The Gaussian projection satisfies the following property:

$$G
ho = \operatorname*{arg\,min}_{\pi\in\mathcal{G}} d_{\mathrm{KL}}(
ho||\pi), \qquad ext{where } \mathcal{G} = ext{Gaussians}$$

• Next step: approximation by a particle system:

$$\varrho_{n+1}^G = GL_n S^J P \varrho_n^G, \qquad S^J \rho := \frac{1}{J} \sum_{j=1}^J \delta_{u^{(j)}}, \qquad u^{(j)} \sim \rho \text{ i.i.d.}$$

What we want to show

In the near-Gaussian setting, the probability measures $\{\rho_n^G\}_{n\in [\![1,N]\!]}$ are close to the true filtering distributions $\{\rho_n\}_{n\in [\![1,N]\!]}$.

Metric on probability measures: weighted total variation

$$d_g(\mu,\nu) = \sup_{|f| \le g} |\mu[f] - \nu[f]|, \quad g: u \to 1 + |u|^2, \quad \rho[f] := \int f d\rho.$$

Assumptions:

• (Near-Gaussian setting) There is $\varepsilon > 0$ such that

 $d_g(G\rho_n,\rho_n) \leq \varepsilon \qquad \forall n \in \llbracket 0,N \rrbracket.$

• (Boundedness) There is $\kappa < \infty$ such that

$$\|\Psi\|_{L^{\infty}(\mathbf{R}^{d})} \vee \|h\|_{L^{\infty}(\mathbf{R}^{d})} \leq \kappa.$$

• (Non-degenerate noise) $\Sigma > 0$ and $\Gamma > 0$.

Strategy of proof

Let \mathcal{P}_R denote the set of probability measures with bounded first and second moments:

$$\mathcal{P}_{R}(\mathbf{R}^{d}) = \left\{ \mu \in \mathcal{P}(\mathbf{R}^{d}) : \max\left\{ |\mathcal{M}(\mu)|, \|\mathcal{C}(\mu)\|, \|\mathcal{C}(\mu)\|^{-1} \right\} \le R \right\}.$$

Auxiliary results:

- The map P is Lipschitz over $\mathcal{P}(\mathbf{R}^d)$ for d_g .
- For any R > 1, the maps L_n and G are Lipschitz over $\mathcal{P}_R(\mathbf{R}^d)$ for d_g .
- Moment bounds: $P\rho_n, P\rho_n^G, L_n P\rho_n, L_n P\rho_n^G$ belong to $\mathcal{P}_{R_*}(\mathbf{R}^d)$ for some R_* .

Main theorem

Under the assumptions,

$$\forall n \in [\![0, N]\!], \qquad d_g(\rho_n^G, \rho_n) \le \varepsilon \left(\frac{\ell^n - 1}{\ell - 1}\right)$$

Proof. Denoting by ℓ the Lipschitz constant of GL_nP over $\mathcal{P}_{R_*}(\mathbf{R}^d)$, we have $d_g(\rho_{n+1}^G, \rho_{n+1}) \leq d_g(\rho_{n+1}^G, G\rho_{n+1}) + d_g(G\rho_{n+1}, \rho_{n+1})$ $= d_g(GL_nP\rho_n^G, GL_nP\rho_n) + d_g(G\rho_{n+1}, \rho_{n+1})$ $\leq \ell d_g(\rho_n^G, \rho_n) + \varepsilon \leq \dots$ Recall the definition

$$P\rho(u) = \int_{\mathbf{R}^d} \frac{\exp\left(-\frac{1}{2} \left|u - \Psi(v)\right|_{\Sigma}^2\right)}{\sqrt{(2\pi)^d \det \Sigma}} \,\rho(v) \,\mathrm{d}v =: \int p(v, u) \,\rho(v) \,\mathrm{d}v.$$

Take any $f \leq g$. Since $\Psi(v)$ is bounded by assumption,

$$b(v) := \int f(u)p(v,u) \, \mathrm{d}u \le \int g(u)p(v,u) \, \mathrm{d}u$$
$$= |\Psi(v)|^2 + \operatorname{tr}(\Sigma) \le \kappa^2 + \operatorname{tr}(\Sigma).$$

Therefore,

$$\begin{aligned} \left| \mathbf{P}\mu[f] - \mathbf{P}\nu[f] \right| &= \left| \int \left(\int f(u)p(v,u) \,\mathrm{d}u \right) \left(\mu(v) - \nu(v) \right) \,\mathrm{d}v \right| \\ &= \left| \mu[b] - \nu[b] \right| \le \left(\kappa^2 + \operatorname{tr}(\Sigma) \right) d_g(\mu,\nu). \end{aligned}$$

^[4] P. REBESCHINI and R. van HANDEL. Can local particle filters beat the curse of dimensionality? Ann. Appl. Probab., 2015.

Auxiliary result 2: the analysis is locally Lipschitz

Recall the definition

$$L_n \mu(u) = \frac{\exp\left(-\frac{1}{2} \left|y_{n+1}^{\dagger} - h(u)\right|_{\Gamma}^2\right) \mu(u)}{\int_{\mathbf{R}^d} \exp\left(-\frac{1}{2} \left|y_{n+1}^{\dagger} - h(U)\right|_{\Gamma}^2\right) \mu(U) \,\mathrm{d}U}$$

Since \boldsymbol{h} is bounded by assumption, there exists \boldsymbol{K} such that

$$\forall u \in \mathbf{R}^d, \qquad K^{-1} \le \phi_n(u) := \exp\left(-\frac{1}{2} \left|y_{n+1}^{\dagger} - h(u)\right|_{\Gamma}^2\right) \le K.$$

For any $f \leq g$, we have

$$\begin{aligned} \left| L_{n}\mu[f] - L_{n}\nu[f] \right| &= \left| \frac{\mu[f\phi_{n}]}{\mu[\phi_{n}]} - \frac{\nu[f\phi_{n}]}{\nu[\phi_{n}]} \right| \\ &\leq \left| \frac{\mu[f\phi_{n}] - \nu[f\phi_{n}]}{\mu[\phi_{n}]} \right| + \left| \frac{\nu[f\phi_{n}](\nu[\phi_{n}] - \mu[\phi_{n}])}{\mu[\phi_{n}]\nu[\phi_{n}]} \right| \\ &\leq K^{2}d_{g}(\mu,\nu) + K^{4}\nu[g] d_{g}(\mu,\nu) = \left(K^{2} + K^{4}\nu[g] \right) d_{g}(\mu,\nu). \end{aligned}$$

This is established in two steps:

• Control $|\mathcal{M}(\mu) - \mathcal{M}(\nu)|$ and $|\mathcal{C}(\mu) - \mathcal{C}(\nu)|$ using $d_g(\mu, \nu)$. For example:

$$|\mathcal{M}(\mu) - \mathcal{M}(\nu)| = \sup_{|\mathbf{a}|=1} \left| \mathbf{a}^{\mathsf{T}} \left(\mathcal{M}(\mu) - \mathcal{M}(\nu) \right) \right|$$
$$= \sup_{|\mathbf{a}|=1} \left| \mu[\mathbf{a}^{\mathsf{T}}u] - \nu[\mathbf{a}^{\mathsf{T}}u] \right| \le d_g(\mu, \nu).$$

\rightsquigarrow The weight in d_g is essential for this step!

Show that, for any two Gaussian measures $\mu = N(\mathbf{m}_1, S_1)$ and $\nu = N(\mathbf{m}_2, S_2)$,

$$d_{g}(\mu,\nu) \leq \sqrt{\left(\mu[g^{2}] + \nu[g^{2}]\right)} \left(3 \left\|S_{2}^{-1}S_{1} - I_{d}\right\|_{F} + \left|\mathbf{m}_{1} - \mathbf{m}_{2}\right|_{S_{2}}\right),$$

This generalizes a similar result for $d_1^{[5]}$.

^[5] L. DEVROYE, A. MEHRABIAN, and T. REDDAD. The total variation distance between high-dimensional Gaussians. arXiv e-prints, 2018.

Extending the error estimate to the particle approximation

Goal: Generalize the approach to the iteration

$$\varrho_{n+1}^G = GL_n S^J P \varrho_n^G, \qquad S^J \mu := \frac{1}{J} \sum_{j=1}^J \delta_{u^{(j)}}, \qquad u^{(j)} \sim \mu \text{ i.i.d.}$$

Since (ϱ_n^G) are random measures, we extend the definition of d_g to the random setting:

$$d_g(\mu,
u) := \sup_{f\leq g} \mathbf{E} \sqrt{\left|\mu[f] -
u[f]\right|^2}.$$

The sampling operator satisfies^[6]

$$d_g(\mu, \mathbf{S}^J \mu) \le \frac{1}{\sqrt{J}} \mathbf{E} \Big(1 + |\mathcal{M}(\mu)|^2 + \operatorname{tr} \big(\mathcal{C}(\mu) \big) \Big)$$

• but the Lipschitz continuity of L_n and G is difficult to show for random measures...

$$\left|L_n\mu[f] - L_n\nu[f]\right| \le \left(K^2 + K^4\nu[g]\right) \left|\mu[\phi_n f] - \nu[\phi_n f]\right|.$$

[6] D. SANZ-ALONSO, A. M. STUART, and A. TAEB. Inverse Problems and Data Assimilation with Connections to Machine Learning. arXiv preprint, 2018. The discrete-time filtering problem

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Rewriting the evolution of the true filtering distribution

Filtering problem

$$\begin{array}{ll} \mbox{Stochastic dynamics:} & u_{n+1} = \Psi(u_n) + \xi_n, & \xi_n \sim \mathsf{N}(0, \varSigma), \\ \mbox{Data model:} & y_{n+1} = h(u_{n+1}) + \eta_{n+1}, & \eta_{n+1} \sim \mathsf{N}(0, \varGamma). \end{array}$$

The true filtering evolves according to

$$\rho_{n+1} = \underline{L}_n \underline{P} \rho_n = \underline{B}^n Q \underline{P} \rho_n.$$

•
$$Q: \mathcal{P}(\mathbf{R}^d) \to \mathcal{P}(\mathbf{R}^d \times \mathbf{R}^K)$$

 $Q\rho(u, y) = \exp\left(-\frac{1}{2}|y - h(u)|_{\Gamma}^2\right)\rho(u).$
• $B^n: \mathcal{P}(\mathbf{R}^d \times \mathbf{R}^K) \to \mathcal{P}(\mathbf{R}^d)$ is the conditioning on the observation $y_{n+1} = y_{n+1}^{\dagger}$

1.

$$\mathbf{B}^n \rho(u) = \frac{\rho(u, y_{n+1}^{\mathsf{T}})}{\int \rho(u, y_{n+1}^{\mathsf{T}}) \, \mathrm{d}u}$$

Schematically,

$$\mathbf{P}(u_n|Y_n) \xrightarrow{P} \mathbf{P}(u_{n+1}|Y_n) \xrightarrow{Q} \mathbf{P}((u_{n+1}, y_{n+1}) \mid Y_n) \xrightarrow{B^n} \mathbf{P}(u_{n+1}|Y_{n+1})$$

Analysis of the ensemble Kalman filter in the near-Gaussian setting

The ensemble Kalman filter from a mean field perspective^[7] (1/2)

One iteration of ensemble Kalman at the mean field level

$$\begin{aligned} \widehat{u}_{n+1} &= \psi(u_n) + \xi_n, & \xi_n \sim \mathsf{N}(0, \Sigma), \\ \widehat{y}_{n+1} &= h(\widehat{u}_{n+1}) + \eta_{n+1}, & \\ u_{n+1} &= \widehat{u}_{n+1} + \mathcal{C}^{uy}(\widehat{\pi}_{n+1}) \mathcal{C}^{yy}(\widehat{\pi}_{n+1})^{-1} (y_{n+1}^{\dagger} - \widehat{y}_{n+1}), & \eta_{n+1} \sim \mathsf{N}(0, \Gamma). \end{aligned}$$

Here $\widehat{\pi}_{n+1} = \operatorname{Law}(\widehat{u}_{n+1}, \widehat{y}_{n+1})$ and, for $\pi \in \mathcal{P}(\mathbf{R}^d \times \mathbf{R}^K)$,

$$\mathcal{C}(\pi) = \begin{pmatrix} \mathcal{C}^{uu}(\pi) & \mathcal{C}^{uy}(\pi) \\ \mathcal{C}^{uy}(\pi)^{\mathsf{T}} & \mathcal{C}^{yy}(\pi) \end{pmatrix}.$$

The third equation may be rewritten

$$u_{n+1} = \underline{T}^n(\widehat{u}_{n+1}, \widehat{y}_{n+1}; \widehat{\pi}_{n+1}),$$

with T^n the following mean-field transport map:

$$T^{n}(\bullet,\bullet;\pi) \colon \mathbf{R}^{d} \times \mathbf{R}^{K} \to \mathbf{R}^{d};$$
$$(u,y) \mapsto u + \mathcal{C}^{uy}(\pi)\mathcal{C}^{yy}(\pi)^{-1}(y_{n+1}^{\dagger} - y)$$

In the following, we use the shorthand notation $T^n_{\sharp}\pi = \left(T^n(\bullet, \bullet; \pi)\right)_{\sharp}\pi$.

^[7] E. CALVELLO, S. REICH, and A. M. STUART. Ensemble Kalman Methods: A Mean Field Perspective. arXiv preprint, 2022.

The ensemble Kalman filter from a mean field perspective

One iteration of ensemble Kalman at the mean field level

$$\begin{aligned} \widehat{u}_{n+1} &= \Psi(u_n) + \xi_n, & \xi_n \sim \mathsf{N}(0, \Sigma), \\ \widehat{y}_{n+1} &= h(\widehat{u}_{n+1}) + \eta_{n+1}, & \\ u_{n+1} &= \widehat{u}_{n+1} + \mathcal{C}^{uy}(\widehat{\pi}_{n+1}) \mathcal{C}^{yy}(\widehat{\pi}_{n+1})^{-1} \big(y_{n+1}^{\dagger} - \widehat{y}_{n+1} \big), & \eta_{n+1} \sim \mathsf{N}(0, \Gamma). \end{aligned}$$

Let $\rho_n^K = \text{Law}(u_n)$. Then

$$\begin{split} \rho_{n+1} &= B^n Q P \rho_n & (\text{True filtering distribution}) \\ \rho_{n+1}^K &= T^n_{\sharp} Q P \rho_{n+1}^K & (\text{Ensemble Kalman filtering distribution}) \end{split}$$

Schematically, for mean field ensemble Kalman:

$$\operatorname{Law}(u_n) \xrightarrow{P} \operatorname{Law}(\widehat{u}_{n+1}) \xrightarrow{Q} \operatorname{Law}((\widehat{u}_{n+1}, \widehat{y}_{n+1})) \xrightarrow{T_{\sharp}^n} \operatorname{Law}(u_{n+1})$$

Key result for the analysis: For any Gaussian ρ

$$T^n_{\sharp}\rho = B^n\rho.$$

 \rightsquigarrow as expected, mean field ensemble Kalman is exact in the Gaussian setting.

What we want to show

In the near-Gaussian setting, the probability measures $\{\rho_n^K\}_{n \in [\![1,N]\!]}$ are close to the true filtering distributions $\{\rho_n\}_{n \in [\![1,N]\!]}$.

Assumptions:

• (Near-Gaussian setting) There is $\varepsilon > 0$ such that

 $d_g(QP\rho_n, GQP\rho_n) \le \varepsilon \qquad \forall n \in [[0, N]].$

• (Boundedness) There is $\kappa < \infty$ such that

$$\|\Psi\|_{L^{\infty}(\mathbf{R}^{d})} \vee \|h\|_{L^{\infty}(\mathbf{R}^{d})} \leq \kappa.$$

- (Lipschitz continuity) The map h is globally Lipschitz continuous.
- (Non-degenerate noise) $\Sigma > 0$ and $\Gamma > 0$.

Main theorem: There exists C_N independent of ε such that

$$\forall n \in [\![0,N]\!], \qquad d_g(\rho_n^K,\rho_n) \le C_N \varepsilon.$$

Strategy of proof

Auxiliary results:

- The maps P, L_n and G are Lipschitz on $\mathcal{P}_R(\mathbf{R}^d)$ with constant $\ell(R)$.
- The maps B^n and T^n_{\sharp} satisfy: $\forall (\mu, \nu) \in \mathcal{P}(\mathbf{R}^d) \times \mathcal{P}_R(\mathbf{R}^d \times \mathbf{R}^K)$,

$$\begin{aligned} d_g(B^n Q P \mu, B^n \nu) &\leq \ell(R) \, d_g(Q P \mu, \nu) \\ d_g(T^n_{\sharp} Q P \mu, T^n_{\sharp} \nu) &\leq \ell(R) \, d_g(Q P \mu, \nu) \end{aligned}$$

 \rightsquigarrow "Lipschitz" when the first argument is in the range of QP. • Moment bounds: all the appropriate measures are in \mathcal{P}_{R_*} .

The main idea of the proof is to use the triangle inequality. Since

$$d_g(\rho_{n+1}^K, \rho_{n+1}) = d_g\left(T^n_{\sharp}Q^P \rho_n^K, B^n Q^P \rho_n\right)$$

and " $T^n_{\sharp}G = B_nG$ " we have

$$d_g \left(T^n_{\sharp} Q P \rho_n^K, B^n Q P \rho_n \right) \leq d_g \left(T^n_{\sharp} Q P \rho_n^K, T^n_{\sharp} Q P \rho_n \right) + d_g \left(T^n_{\sharp} Q P \rho_n, T^n_{\sharp} G Q P \rho_n \right) + d_g \left(B^n G Q P \rho_n, B^n Q P \rho_n \right) \leq \ell(R_*)^3 d_g(\rho_n^K, \rho_n) + 2\ell(R_*)\varepsilon.$$

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Conclusions and perspectives

In this presentation,

- we confined ourselves to the mean field setting;
- we analysed a simple filter based on Gaussian projections;
- we analysed the ensemble Kalman filter in the near-Gaussian setting.

Perspectives for future work:

- Obtain error estimates with a better scaling with respect to N;
- Obtain error estimates for continuous-time Gaussian filtering methods;
- Derive error bounds for the particle approximations;
- Identify settings in which the "near-Gaussian" assumption is provably correct.

Some references

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Thank you for your attention!